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Ultrathin 2D/2D WO₃/g-C₃N₄ step-scheme H₂-production photocatalyst

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ABSTRACT

The appropriate interfacial contact of heterojunction photocatalysts plays a critical role in transfer/separation of interfacial charge carriers. Design of two-dimensional (2D)/2D surface-to-surface heterojunction is an effective method for improving photocatalytic activity since greater contact area can enhance interfacial charge transfer rate. Herein, ultrathin 2D/2D WO₃/g-C₃N₄ step-like composite heterojunction photocatalysts were fabricated by electrostatic self-assembly of ultrathin tungsten trioxide (WO₃) and graphitic carbon nitride (g-C₃N₄) nanosheets. The ultrathin WO₃ and g-C₃N₄ nanosheets were obtained by electrostatic-assisted ultrasonic exfoliation of bulk WO₃ and a two-step thermal-etching of bulk g-C₃N₄, respectively. The thickness of ultrathin WO₃ and g-C₃N₄ nanosheets are 2.5–3.5 nm, which is equivalent to 5–8 atomic or molecular layer thickness. This ultrathin layered heterojunction structure can enhance surface photocatalytic rate because photogenerated electrons and holes at heterogeneous interface more easily transfer to surface of photocatalysts. Therefore, the obtained ultrathin 2D/2D WO₃/g-C₃N₄ step-scheme (S-scheme) heterojunction photocatalysts exhibited better H₂-production activity than pure g-C₃N₄ and WO₃ with the same loading amount of Pt as cocatalyst. The mechanism and driving force of charge transfer and separation in S-scheme heterojunction photocatalysts are investigated and discussed. This investigation will provide new insight about designing and constructing novel S-scheme heterojunction photocatalysts.

1. Introduction

Photocatalytic hydrogen (H_2) generation has been considered as one of the most promising routes to convert solar energy into available chemical energy [1–7]. However, single-semiconductor photocatalysts cannot reach high photocatalytic activity due to their high recombination probability of photogenerated charge carriers and their limited redox potential. Constructing suitable heterojunction systems is an effective way to solve this issue [8–15]. Generally, the design of high efficiency H_2 -generation heterojunction photocatalysts mainly focuses on two key points. One is the suitable band staggered arrangement of two semiconductors, and the other is the ideal interfaces for charge transfer/separation between them [16].

Since Wang et al. firstly used graphitic carbon nitride (g- C_3N_4) for H_2 generation by photocatalytic water splitting [17], g- C_3N_4 attracted much attention due to its narrow bandgap with visible-light response, high conduction band (CB) position with high reduction ability, simple syntheses, and special two-dimensional (2D) layered structure

[5,8,11,18-21]. However, the photocatalytic performance of pure g-C₃N₄ is still far from the practical requirement [22]. A variety of semiconductors with staggered band structure relative to g-C₃N₄ have been used to couple with g-C₃N₄ for constructing g-C₃N₄-based heterogeneous photocatalysts, such as TiO₂/g-C₃N₄ [18,23-26], ZnO/g- C_3N_4 [27–30], WO_3/g - C_3N_4 [31–35], CdS/g- C_3N_4 [36–38], $ZnIn_2S_4/g$ -C₃N₄ [39], BiOI/g-C₃N₄ [40]. At the interfaces of these type-II heterojunctions, the photogenerated electrons will transfer from the semiconductor with higher conduction band (CB) to that with lower CB, while the photogenerated holes will transfer from the semiconductor with lower valence band (VB) to that with higher VB. As shown in Fig. 1a, the interfacial charge transfer in conventional type-II heterojunctions greatly reduces the redox abilities of photogenerated electrons and holes. This kind of charge transfer mechanism thermodynamically is not beneficial for the occurrence of photocatalytic oxidation and reduction reactions [8]. Moreover, the photogenerated electrons in CB of photocatalyst II (PC II) are difficult to transfer to CB of photocatalyst I (PC I) due to their strong Coulomb electrostatic

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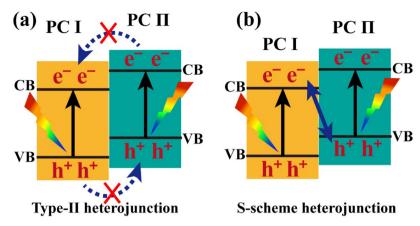


Fig. 1. The schematic diagrams of charge transfer in conventional type-II heterojunction (a) and S-scheme heterojunction (b).

repulsive force [41,42]. Also, the Coulombic attraction force between the electrons (in CB of PC II) and holes (in VB of PC II) also hinders this electron transfer. Similarly, the holes in VB of PC I also cannot easily transfer to VB of PC II. Therefore, the charge transfer mechanism of conventional type-II heterojunction is with obvious limitations from viewpoint of kinetics.

To overcome the shortcomings of conventional type-II heterojunction, meanwhile, to combine our previous work on Z-scheme photocatalyst [41–50], a new step-scheme (S-scheme) heterojunction concept is proposed and illustrated in Fig. 1b. This S-scheme heterojunction photocatalyst is mainly composed of two n-type semiconductor photocatalysts, and PC I and PC II represent oxidation photocatalyst and reduction photocatalyst, respectively. The transfer of photogenerated electrons in S-scheme heterojunction is more like "step" (macroscopic viewpoint) or "N" (microscopic viewpoint) type. The electrons and holes are separated in space, locating in CB of PC II and VB of PC I, respectively. The driving force of charge carrier transfer is mainly from the internal electric field between PC I and PC II. Usually, the oxidation photocatalyst PC I has greater work function and lower Fermi level; contrarily, the reduction photocatalyst PC II has smaller work function and higher Fermi level. When PC I and PC II contact, the electrons in PC II will transfer to PC I across their interface. An internal electric field will be built and its direction is from PC II to PC I. In S-scheme heterojunction, comparatively useless electrons in the CB of PC I and comparatively useless holes in the VB of PC II are recombined and eliminated at the interface. Contrarily, useful holes in VB of PC I and useful electrons in CB of PC II are kept due to the presence of the internal electric field. Finally, photocatalytic oxidation and reduction reactions are initiated by holes in the VB of PC I and electrons in the CB of PC II, respectively.

Herein, ultrathin 2D tungsten trioxide (WO $_3$) and g-C $_3$ N $_4$ nanosheets are first prepared. Then, 2D/2D WO $_3$ /g-C $_3$ N $_4$ S-scheme heterogeneous composite photocatalysts are constructed with an electrostatic self-assembly method. The prepared samples are carefully characterized by X-ray diffraction (XRD), transmission electron microscopy (TEM), atomic force microscopy (AFM), X-ray photoelectron spectroscopy (XPS), Raman spectroscopy and Zeta potential measurement, etc. The H $_2$ -production performance of the samples is evaluated by water splitting using lactic acid as a sacrificial agent. The S-scheme heterojunction photocatalytic mechanism in 2D/2D WO $_3$ /g-C $_3$ N $_4$ composites is in depth analyzed and discussed.

2. Experimental section

2.1. Preparation of bulk WO3

Firstly, 500 mg of $\rm Na_2WO_4.2H_2O$ was dissolved in 200 mL of $\rm HNO_3$ solution (4.8 M). The mixture solution was stirred for 36 h. Then, the

yellow precipitate (WO $_3$ 2H $_2$ O) was collected by centrifugation, and washed with water until the neutral pH. The collected WO $_3$ 2H $_2$ O was dried for 12 h. The bulk WO $_3$ was obtained by directly calcining the dried WO $_3$ 2H $_2$ O powder with a 5 °C/min heating rate and keeping at 500 °C for 3 h.

2.2. Preparation of ultrathin 2D WO₃ nanosheets

The ultrathin WO₃ nanosheets were obtained by electrostatic-assisted ultrasonic exfoliation of bulk WO3, with bovine serum albumin (BSA) used as an exfoliating agent. The abundant surface -NH₂ groups of BSA exhibit strong electrostatic binding with WO3 under acidic condition. This strong electrostatic force can facilitate the dissociation of WO₃ nanosheets from the surface of bulk WO₃, and greatly improve the dispersity of WO₃ nanosheets in solution [51]. In a typical experiment, 10 mg of BSA was dissolved in 100 mL of H₂O, and then the pH of the mixture solution was adjusted to 4 with 1 M HNO₃. 50 mg of bulk WO₃ powder was dispersed in the above solution and vigorously sonicated for 3 h. The resulting milky suspension was subjected to centrifugation at 4500 rpm for 30 min. After removing the supernatant BSA solution, the precipitant was re-dispersed in 100 mL of H2O with pH = 4 and vigorously sonicated for another 1 h. Finally, the white milky WO3 suspension was obtained and stored prior to use. The concentration of WO₃ nanosheets in the suspension was 0.5 mg/mL.

2.3. Preparation of ultrathin 2D g-C₃N₄ nanosheets

g- C_3N_4 nanosheets were prepared by a two-step calcination thermaletching method using urea as precursor [50]. The detailed methods are shown in the Supporting information.

2.4. Preparation of ultrathin 2D/2D WO₃/g-C₃N₄ composite

 $WO_3/g\text{-}C_3N_4$ composites were prepared by electrostatic self-assembly of WO_3 nanosheets and $g\text{-}C_3N_4$ nanosheets. Their preparation method is provided in the Supporting information.

The detailed methods for characterization, photocatalytic H_2 generation test and computational calculations are given in the Supporting information.

3. Results and discussion

3.1. Zeta potential

Fig. 2a exhibits the zeta potentials of bulk WO₃, WO₃ nanosheets and g-C₃N₄ nanosheets at pH = 4. The bulk WO₃ shows a negative zeta potential of -9.7 mV at pH = 4, whereas the WO₃ nanosheets exhibits a more negative zeta potential of -22.8 mV at the same pH. The higher

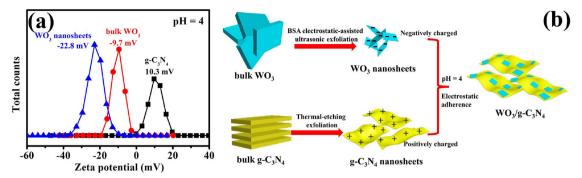


Fig. 2. (a) Zeta potentials of bulk WO_3 , WO_3 nanosheets and $g-C_3N_4$ at pH=4. (b) The formation schematic diagram of 2D/2D $WO_3/g-C_3N_4$ heterojunctions by Coulomb electrostatic interaction.

zeta potential value indicates better dispersity of WO $_3$ nanosheets than bulk WO $_3$. The exfoliation exposes more surface groups and improves the dispersity of WO $_3$ nanosheets. The g-C $_3$ N $_4$ nanosheets exhibit a positive zeta potential of 10.3 mV at pH = 4. The opposite zeta potentials result in strong electrostatic attraction between WO $_3$ and g-C $_3$ N $_4$ nanosheets [52], which is beneficial for the charge transfer between them. Thus, stable 2D/2D WO $_3$ /g-C $_3$ N $_4$ heterojunctions are obtained by the Coulomb electrostatic interaction. The fabrication schematic diagram of the 2D/2D WO $_3$ /g-C $_3$ N $_4$ heterojunctions is shown in Fig. 2b.

3.2. Morphological characteristics

TEM analysis is a powerful tool to study the ultrathin nanosheets. As shown in Fig. 3a, clear nanosheet structure is observed for the WO₃ nanosheet sample, with the lateral size of WO₃ nanosheets measured at about 50×50 nm. The low contrast of TEM image indicates ultrathin thickness of WO₃ nanosheets, which is further investigated by AFM (Fig. 4a and b). Clearly, the lateral size of WO₃ nanosheets is between 50 and 100 nm, and the thickness of WO₃ nanosheets is 3-4 nm, which directly prove the presence of ultrathin WO₃ nanosheets (with 5-8 atom or molecular layers). The HRTEM image of WO3 nanosheets (corresponding to the part in the red circle in Fig. 3a) is shown in Fig. 3b, and clear lattice fringes are observed. Combined with the electronic diffraction pattern (inset of Fig. 3b), the WO₃ nanosheets are single crystal nanosheets with good crystallinity. The measured fringe spacings are 0.377 and 0.365 nm, which can be assigned to the (020) and (200) facets of monoclinic WO3, respectively. Fig. 3c presents TEM image of g-C₃N₄. Flexible nanosheets with large lateral size can be observed, which is typical characteristics of g-C₃N₄ nanosheets. The HRTEM image (Fig. 3d, corresponding to the part in the red circle in Fig. 3c) indicates that the thickness of g- C_3N_4 nanosheets is also very small. The actual thickness of g-C₃N₄ nanosheets was also measured by the AFM. As shown in Fig. 4c and 4d, the g-C₃N₄ nanosheets show microscale lateral size and the thickness of g-C₃N₄ is about 2.5 nm. The interlayer spacing of g-C₃N₄ is 0.36 nm, and thus the ultrathin g-C₃N₄ nanosheets are about seven molecular layers in thickness. In the TEM image of 15% WO₃/g-C₃N₄ sample (Fig. 3e), the circled parts with higher contrast can be assigned to the WO₃ nanosheets, whose sizes are consistent with the pure WO₃ nanosheets shown in Figs. 3a and 4 a. For further observation, HRTEM image (Fig. 3f) exhibits the clear contact interface of WO₃ and g-C₃N₄. Clear lattice fringes in the upper part of the image can be ascribed to WO3. The interfacial contact of WO3 and g-C3N4 is very close, which can be attributed to the strong interfacial electrostatic adherence force. This intimate interfacial contact is beneficial for the interfacial charge transfer between WO3 and g-C3N4.

3.3. Phase structure and surface chemical state

Fig. S1 (Supporting information) compares XRD patterns of

exfoliated WO $_3$ nanosheets, bulk WO $_3$ and WO $_3$ 2H $_2$ O. All the typical peaks of WO $_3$ 2H $_2$ O sample can be indexed to the corresponding reflections of monoclinic WO $_3$ 2H $_2$ O (JCPDS PDF No. 40–0693). XRD pattern of WO $_3$ nanosheets is similar to that of bulk WO $_3$, which means that exfoliation treatment does not change the main crystal structure of WO $_3$. All the peaks are ascribed to the monoclinic WO $_3$ (JCPDS PDF No. 83-0950) [53], and no other impurity peak can be observed, implying the high purity of WO $_3$ nanosheets products. It is noteworthy that XRD peaks of exfoliated WO $_3$ nanosheets become much weaker and broader than that of bulk WO $_3$, which can be ascribed to the thinner nanosheet structure of exfoliated WO $_3$ nanosheets. This indicates the successful exfoliation of bulk WO $_3$ through the BSA electrostatic-assisted ultrasonic exfoliation method.

Fig. 5a compares XRD patterns of g-C₃N₄ nanosheets, WO₃ nanosheets and WO₃/g-C₃N₄ composites with different WO₃ percentages. The XRD pattern of pure g-C₃N₄ exhibits a typical peak located at 27.5°, ascribed to the characteristic interlayer stacking of aromatic system [22,54,55]. Another characteristic peak located at about 13° (in-plane structural packing motif of tri-s-triazine units) is too weak to be directly observed. Moreover, a clear broad peak located at about 22° can be observed, which is a characteristic peak of amorphous matter. This observation indicates the successful exfoliation of g-C₃N₄ during the second thermal treatment, and ultrathin g-C₃N₄ nanosheets are obtained [56]. In the case of WO₃/g-C₃N₄ composites, the XRD patterns are seen as the combination of XRD patterns of WO₃ nanosheets and g-C₃N₄ nanosheets. With increasing contents of WO₃ nanosheets, the XRD peak intensities of WO₃ become stronger. These results demonstrate the good integration of WO3 nanosheets and g-C3N4 nanosheets in the WO₃/g-C₃N₄ composites.

Fig. 5b compares FTIR spectra of g-C₃N₄, bulk WO₃, WO₃ nanosheets and 15%WO₃/g-C₃N₄. No obvious difference is observed in the spectra of bulk WO₃ and WO₃ nanosheets, and the broad peak located at 450–900 cm⁻¹ is assigned to the stretching vibration of W–O–W [57]. As for pure g-C₃N₄ nanosheets, the absorption bands located at $3000-3500 \, \mathrm{cm^{-1}}$ are assigned to the amino and surface hydroxyl groups. The bands at $1200-1600 \, \mathrm{cm^{-1}}$ are the characteristic of C–N bonds in the tri-s-triazine units of g-C₃N₄. A sharp band at $810 \, \mathrm{cm^{-1}}$ is from the characteristic breathing-vibration of tri-s-triazine units [22,58]. The $15\% WO_3/g$ -C₃N₄ exhibits a similar spectrum with pure g-C₃N₄. Closer examination on the region of $900-1100 \, \mathrm{cm^{-1}}$ (Fig. 5c) reveals a small broad band assigned to WO₃. These results indicate successful combination of WO₃ and g-C₃N₄ in the $15\% WO_3/g$ -C₃N₄.

Fig. 5d shows Raman spectra of g- C_3N_4 , bulk WO_3 , WO_3 nanosheets and $15\%WO_3/g$ - C_3N_4 . Firstly, the WO_3 nanosheets exhibit a similar spectrum to that of its bulk counterpart. The Raman peaks located at 275 and $328\,\mathrm{cm}^{-1}$ are ascribed to the bending vibration mode of O–W–O, while those located at 712 and $810\,\mathrm{cm}^{-1}$ are due to the stretching vibration mode of W–O [53,59]. It is worth noting that Raman peaks of WO_3 nanosheets become broader and weaker as compared to bulk WO_3 , which is due to the phonon softening and

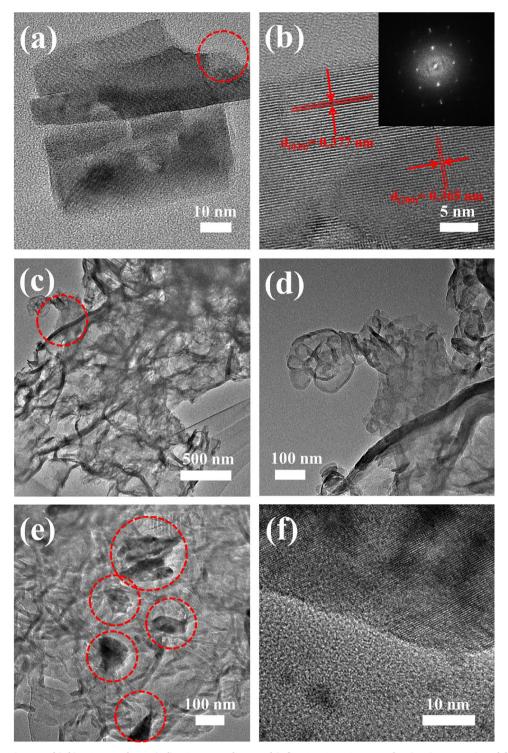


Fig. 3. TEM and HRTEM images of (a,b) WO₃ nanosheets, (c,d) g-C₃N₄ nanosheets and (e,f) 15%WO₃/g-C₃N₄ samples. (For interpretation of the references to colour in the text, the reader is referred to the web version of this article.)

enhanced electron–phonon coupling in the few-layered region [53]. Raman spectra of g-C₃N₄ nanosheets and 15%WO₃/g-C₃N₄ show no obvious peaks, which is attributed to the strong interference by fluorescence effect of g-C₃N₄. The enlarged range within the 750–900 cm $^{-1}$ range are shown in Fig. 5e, in which two small Raman peaks at 712 and 810 cm $^{-1}$ can be observed for 15%WO₃/g-C₃N₄. This result further confirms successful combination of WO₃ and g-C₃N₄ in 15%WO₃/g-C₃N₄.

3.4. XPS analysis

The surface chemical states were further investigated by XPS. All binding energies were calibrated by C 1 s binding energy at 284.8 eV. XPS spectra of bulk WO $_3$ and WO $_3$ nanosheets are presented in Fig. S2 (Supporting information). As shown in Fig. S2a, compared with bulk WO $_3$, two new C 1 s peaks located at 286.5 and 288.0 eV can be observed in the spectrum of WO $_3$ nanosheets, which can be assigned to the C–N and – COOH bands of the surface BSA residual. Fig. S2b exhibits

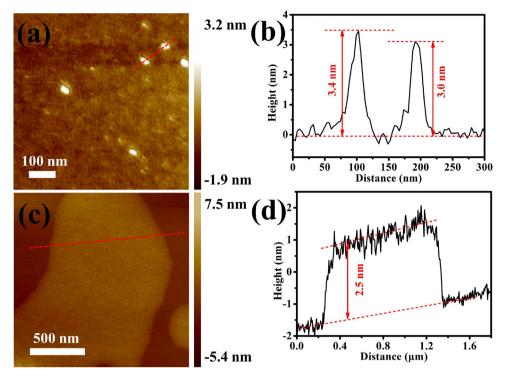


Fig. 4. AFM images of (a,b) WO₃ nanosheets and (c,d) g-C₃N₄ nanosheets.

the comparison of W 4f spectra of bulk WO $_3$ and WO $_3$ nanosheets. Two peaks at 35.9 and 38.1 eV are ascribed to the W 4f $_{7/2}$ and W 4f $_{5/2}$ of W⁶⁺, respectively. A small broad peak at about 41.5 eV is assigned to the W 5p $_{3/2}$. For the WO $_3$ nanosheets, another two peaks at 35.0 and 37.2 eV were observed, which can be attributed to the W 4f $_{7/2}$ and W 4f $_{5/2}$ of W⁵⁺. The W⁵⁺ is derived from the unsaturated W–O bonds on the surface of WO $_3$ nanosheets. The presence of W⁵⁺ indicates

exfoliation of bulk WO_3 into ultrathin WO_3 nanosheets. Fig. S2c shows the O 1s spectra of bulk WO_3 and WO_3 nanosheets. The peak at 530.5 eV is from the lattice oxygen in bulk WO_3 . The WO_3 nanosheets exhibit two peaks at 529.7 and 531.0 eV. The former is assigned to the lattice oxygen at the surface of WO_3 nanosheets, while the latter is from the surface adsorbed oxygen species.

Fig. 6a exhibits C 1s spectra of g-C₃N₄, $15\%WO_3/g$ -C₃N₄ and WO_3

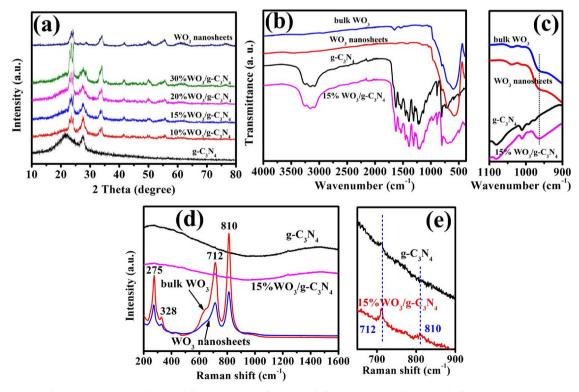


Fig. 5. (a) XRD patterns for g- C_3N_4 , WO_3 nanosheets and their composites. (b) FTIR and (d) Raman spectra for g- C_3N_4 , bulk WO_3 , WO_3 nanosheets and 15% WO_3 /g- C_3N_4 . (c) and (e) are the enlarged parts of (b) and (d), respectively.

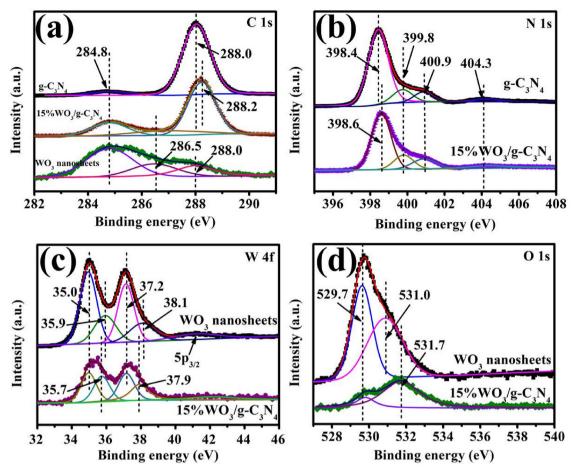


Fig. 6. Comparison of (a) C 1 s, (b) N 1 s, (c) W 4f and (d) O 1s XPS spectra of g-C₃N₄, 15%WO₃/g-C₃N₄ and WO₃ nanosheets.

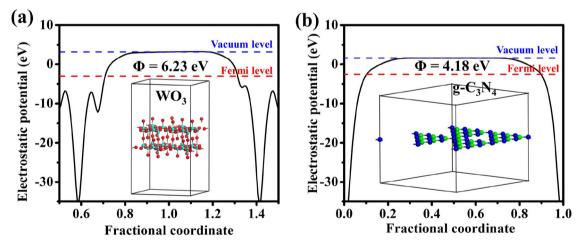


Fig. 7. Electrostatic potentials of (a) WO₃ (001) surface and (b) g-C₃N₄ (001) surface. Insets show the structural models of the materials for DFT calculation.

nanosheets. The XPS peaks at 288.0 and 288.2 eV of g-C₃N₄ and 15% WO₃/g-C₃N₄ are assigned to the N=C-N₂ bonds of g-C₃N₄ components [60]. The slight positive shift of C 1 s peak (by 0.2 eV) is caused by the charge transfer at the interface of g-C₃N₄ and WO₃. The N 1 s spectra of g-C₃N₄ and 15%WO₃/g-C₃N₄ are shown in Fig. 6b. The peak at 399.8 eV is assigned to three-coordinate N-(C)₃, the peak at 400.9 eV to the surface amino group (-NH₂), the peak at 404.3 eV to the π -excitation of C-N heterocycles, and the peaks at 398.4 and 398.6 eV to the two-coordinate N species (C = N-C) [22,60]. The slight positive shift of N 1 s peak (by 0.2 eV) is also attributed to the interfacial charge transfer of g-C₃N₄ and WO₃. To further investigate the charge transfer

between g-C₃N₄ and WO₃ in 15%WO₃/g-C₃N₄, the comparison of W 4f spectra of WO₃ nanosheets and 15%WO₃/g-C₃N₄ is presented in Fig. 6c. The binding energies of W⁶⁺ at 35.9 and 38.1 eV in WO₃ nanosheets negatively shift to 35.7 and 37.9 eV in 15%WO₃/g-C₃N₄. The positive shifts of C 1 s and N 1 s, and the negative shift of W⁶⁺ indicate the charge transfer from g-C₃N₄ to WO₃ at the interfaces [60]. It is noted that, in the O 1 s spectra of WO₃ nanosheets and 15%WO₃/g-C₃N₄ (Fig. 6d), the binding energy of lattice oxygen at 529.7 eV shows no obvious shift, while that of the surface adsorbed oxygen species positively shifts from the 531.0 eV of WO₃ nanosheets to 531.7 eV of 15% WO₃/g-C₃N₄, which is ascribed to the oxidation of adsorbed BSA during

irradiation.

3.5. Density functional theory computational calculations

Density functional theory (DFT) computational calculations were used to further investigate the interfacial charge transfer between WO₃ and g-C₃N₄ [61]. As shown in Fig. S3 (Supporting information), the calculated band gap energies of WO₃ and g-C₃N₄ are 2.77 and 2.68 eV, respectively. Fig. 7 exhibits the electrostatic potentials of WO₃ (001) surface and g-C₃N₄ (001) surface. The work functions of WO₃ (001) surface and g-C₃N₄ (001) surface are calculated to be 6.23 and 4.18 eV, respectively. The difference of work function indicates the presence of charge transfer at the interface of WO3 and g-C3N4. The greater work function of WO3 will lead to the charge transfer from g-C3N4 to WO3 untill the Fermi level equilibrium. This charge transfer will result in the production of interfacial bulid-in electric field. At the interface, the surface of g-C₃N₄ is positively charged, while the surface of WO₃ is negatively charged. This result is consistent with the charge transfer form g-C₃N₄ to WO₃ as revealed by XPS results. Due to the thin layered structures of g-C₃N₄ and WO₃ and their strong Coulomb force, the intensity of bulid-in electric field is greatly enhanced, which is more beneficial for the transfer and separation of interfacial charge carriers than their bulk counterparts.

3.6. Photocatalytic hydrogen generation

Fig. 8a exhibits comparison of photocatalytic H_2 -generation rates of the as-prepared samples using lactic acid as a hole-sacrificial agent. WO₃ nanosheets have no obvious H_2 generation activity observed due to their low CB position. Pure $g\text{-}G_3N_4$ nanosheets show obvious H_2 -generation activity due to their high CB position and photogenerated electrons with enough reduction ability. In the presence of a small amount of WO₃ nanosheets, the H_2 -production activity of WO₃/ $g\text{-}G_3N_4$ composite is obviously enhanced. Especially, the $15\%WO_3/g\text{-}G_3N_4$ composite sample exhibits the highest H_2 -production activity (982 μ mol/h/g), and its activity is about 1.7 times higher than pure $g\text{-}G_3N_4$ nanosheets. Fig. 8b shows that the $15\%WO_3/g\text{-}G_3N_4$ sample exhibits good photocatalytic H_2 -production stability, with no obvious hydrogen activity decrease observed after four cycles.

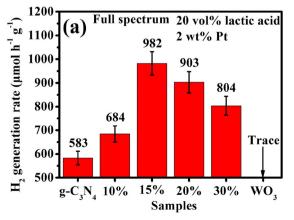
3.7. Photocatalytic mechanism

Fig. 9a exhibits the Mott-Schottky plots of g- C_3N_4 and WO_3 nanosheets. Firstly, both g- C_3N_4 and WO_3 nanosheets show positive slopes, indicating the n-type semiconductor characteristic of g- C_3N_4 and WO_3 nanosheets. The flat band position of n-type semiconductors is

closer to their CB position. Thus, the extrapolated CB positions of g-C $_3$ N $_4$ and WO $_3$ nanosheets are -1.25 and -0.34 V, respectively (vs Ag/AgCl, pH = 7). These two potentials are converted to -0.61 and +0.30 V (vs NHE, pH = 0). Combined with the bandgap of g-C $_3$ N $_4$ (2.75 eV) and WO $_3$ (2.68 eV) nanosheets, their band structures are shown in Fig. 9b.

Fig. 9c shows electron paramagnetic resonance (EPR) spectra of DMPO-'OH measured in aqueous suspensions of the materials. The DMPO-OH signal of WO3 nanosheets and 15%WO3/g-C3N4 can be observed, while no obvious DMPO-OH signal is observed for pure g-C₃N₄. The absence of DMPO-OH signal in g-C₃N₄ is due to the weak oxidation potential of photogenerated holes in g-C₃N₄ (2.14 V, vs NHE at pH = 0). The observation of DMPO- $^{\circ}$ OH signal in the 15%WO₃/g-C₃N₄ composite suggests that the photogenerated holes still stay in the VB of WO₃ and do not transfer to the VB of g-C₃N₄. In Fig. 9d, strong DMPO-O2 signals are observed for g-C3N4 and 15%WO3/g-C3N4 composite samples in methanolic suspension, whereas very weak DMPO-'O₂ signal is observed for WO₃ nanosheets. The results indicate that the photogenerated electrons in g-C₃N₄ and 15%WO₃/g-C₃N₄ composite samples have enough reduction ability to reduce O2 to form superoxide radical anions ('O₂-'). The above EPR results indicate that the photogenerated electrons and holes respectively are present in the CB of g-C₃N₄ and VB of WO₃, and the charge transfer does not follow the conventional type II heterojunction mechanism. Contrarily, the above suggested S-scheme heterojunction mechanism can better explain the enhancement of photocatalytic H₂-production activity in the 2D/2D WO₃/g-C₃N₄ composite photocatalysts.

Fig. 10 shows charge transfer mechanism of S-scheme heterojunction between WO3 and g-C3N4. Usually, g-C3N4 is a reduction-type photocatalyst with smaller work function (4.18 eV) and higher Fermi level. Contrarily, WO₃ is an oxidation-type photocatalyst with larger work function (6.23 eV) and lower Fermi level (Fig. 10a). When g-C₃N₄ and WO₃ are in close contact, the electrons in g-C₃N₄ spontaneously transfer to WO3 across their interface until their Fermi levels are the same (Fig. 10b). Thus, g-C₃N₄ loses electrons and is positively charged, while the WO₃ get electrons and is negatively charged at the interface. Naturally, an internal electric field is produced at the interface. Simultaneously, band edge of g-C₃N₄ bends upward due to the loss of electrons, but band edge of WO3 bends downward due to the accumulation of electrons. Under light irradiation, the electrons are excited from VB to CB of both WO₃ and g-C₃N₄. The internal electric field, band edge bending and Coulomb interaction accelerate the recombination of some electrons (from CB of WO₃) and holes (from VB of g-C₃N₄), and, at the same time, prevent the recombination of some electrons (from CB of g-C₃N₄) and holes (from VB of WO₃) (Fig. 10c). This S-scheme heterojunction mechanism will eliminate the relatively useless electrons



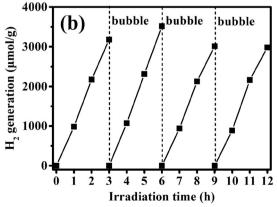


Fig. 8. (a) Comparison of H_2 -generation activities of as-prepared samples under xenon lamp irradiation (full spectrum), with photo-deposition of 2% Pt as cocatalyst (The symbols 10%, 15%, 20% and 30% on the x-axis represent the $10\%WO_3/g$ - C_3N_4 , $15\%WO_3/g$ - C_3N_4 , $20\%WO_3/g$ - C_3N_4 and $30\%WO_3/g$ - C_3N_4 samples, respectively). (b) Time courses of photocatalytic H_2 -generation of $15\%WO_3/g$ - C_3N_4 .

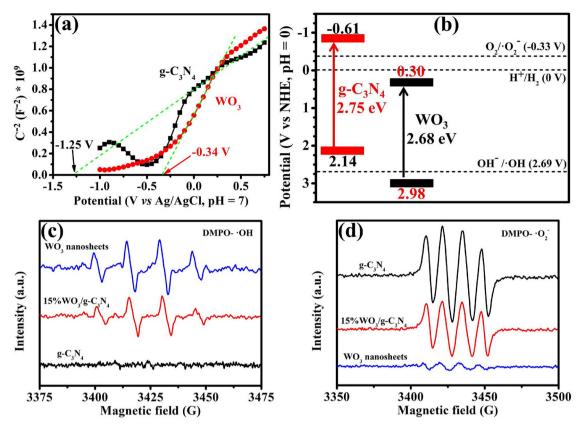


Fig. 9. (a) Mott-Schottky plots and (b) band structures of g- C_3N_4 and WO_3 nanosheets. EPR spectra of (c) DMPO-•OH in aqueous and (d) DMPO-• O_2 in methanol dispersion in the presence of g- C_3N_4 , WO_3 nanosheets and $15\%WO_3/g$ - C_3N_4 .

(from CB of WO₃) and holes (from VB of g-C₃N₄), but hold the useful electrons (from CB of g-C₃N₄) and holes (from VB of WO₃). This charge carrier transfer process endows the 2D/2D WO₃/g-C₃N₄ composite heterojunction with supreme redox capacity, thus providing strong driving force for running photocatalytic water splitting reaction.

4. Conclusions

In summary, ultrathin 2D/2D WO $_3$ /g-C $_3$ N $_4$ S-scheme heterojunction photocatalysts were successfully prepared by an electrostatic self-assembly method using ultrathin WO $_3$ and g-C $_3$ N $_4$ nanosheets as precursors. The H $_2$ -production activity of the WO $_3$ /g-C $_3$ N $_4$ composite

samples was greatly enhanced, which is about 1.7 times higher than pure g- $\rm C_3N_4$. The enhanced photocatalytic performance is due to the formation of step-scheme heterojunction, which inhibits the recombination of useful electrons and holes, but expedites the recombination of relatively useless electrons and holes. XPS, work function and EPR results further confirm the presence of step heterojunction mechanism and its correctness. This work will provide new insight for design and fabrication of novel heterojunction photocatalyst. The prepared ultrathin 2D/2D WO_3/g-C_3N_4 composite can also find potential application in electrocatalysis, catalysis, solar cell and adsorption.

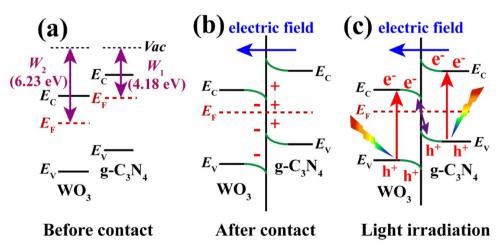


Fig. 10. (a) The work functions of g- G_3N_4 and WO_3 before contact. (b) The internal electric field and band edge bending at the interface of WO_3/g - G_3N_4 after contact. (c) The S-scheme charge transfer mechanism between WO_3 and g- G_3N_4 under light irradiation.

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Appendix A. Supplementary data

Supplementary material related to this article can be found, in the online version, at doi:https://doi.org/10.1016/j.apcatb.2018.11.011.

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